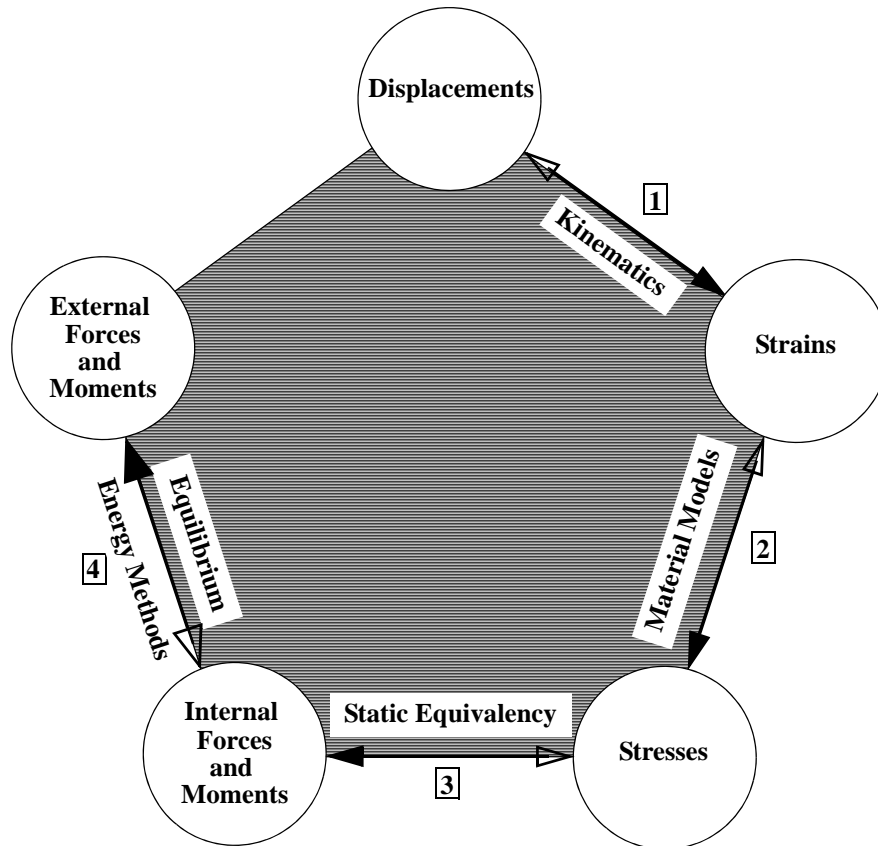


Energy Methods

- Minimum-energy principles are an alternative to statement of equilibrium equations.

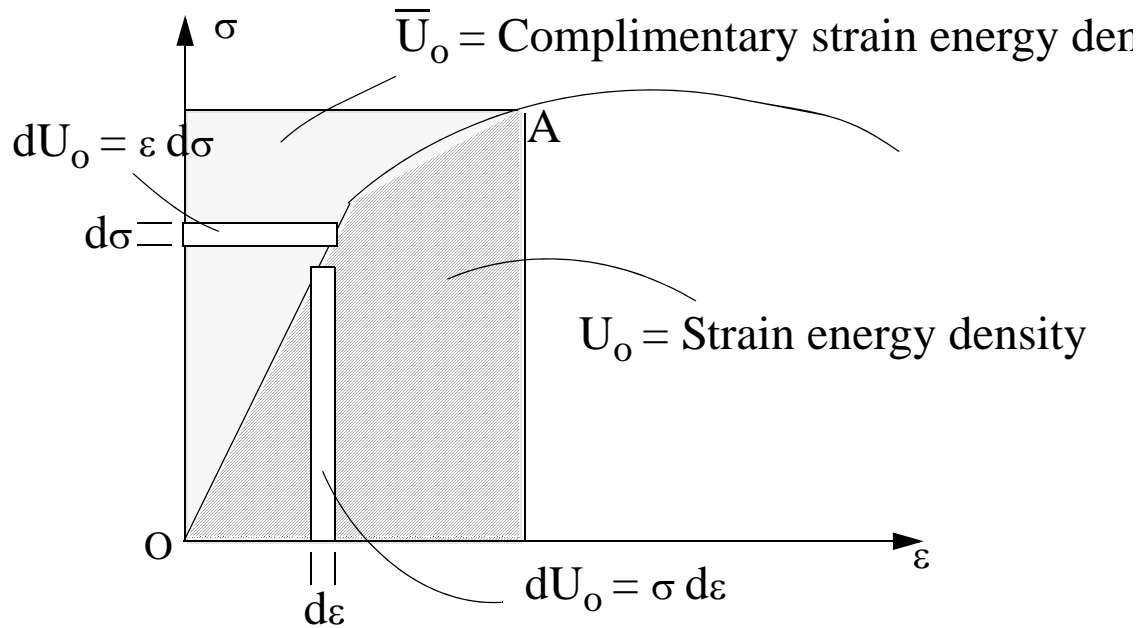


The learning objectives in this chapter are:

- Understand the perspective and concepts in energy methods.
- Learn the use of dummy unit load method and Castigliano's theorem for calculating displacements in statically determinate and indeterminate structures.

Strain Energy

- The energy stored in a body due to deformation is called the *strain energy*.
- The strain energy per unit volume is called the *strain energy density* and is the area underneath the stress-strain curve up to the point of deformation.



Strain Energy:
$$U = \int_V U_o dV [$$

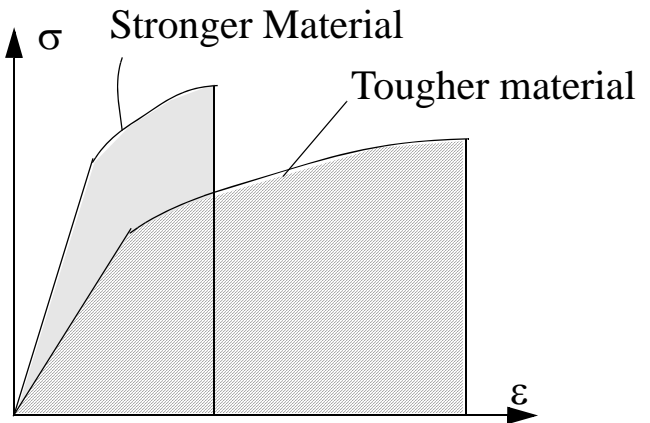
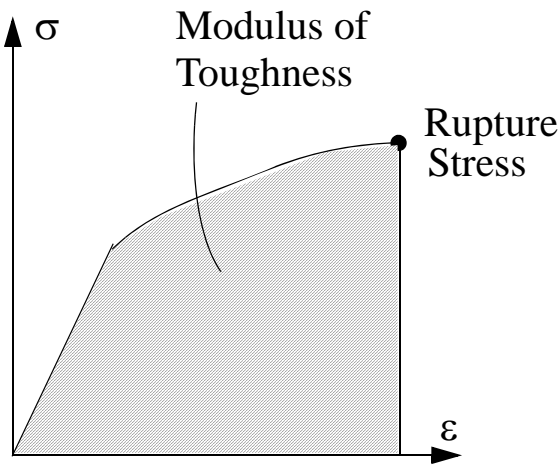
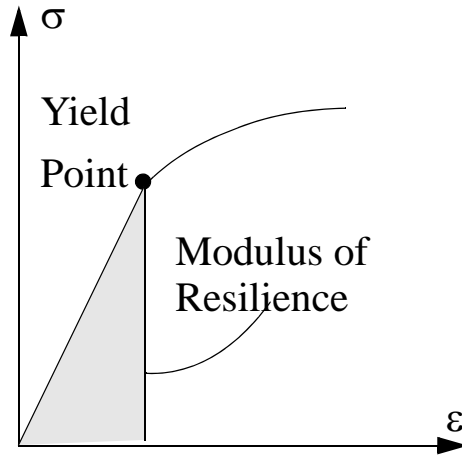
Strain Energy Density:
$$U_o = \int_0^\epsilon \sigma d\epsilon$$

Units: $\text{N}\cdot\text{m} / \text{m}^3, \text{Joules} / \text{m}^3, \text{in}\cdot\text{lbs} / \text{in}^3, \text{ or } \text{ft}\cdot\text{lb}/\text{ft}^3$

Complimentary Strain Energy Density:
$$\bar{U}_o = \int_0^\sigma \epsilon d\sigma$$

- The strain energy density at the yield point is called *Modulus of Resilience*.

- The strain energy density at rupture is called *Modulus of Toughness*.



Linear Strain Energy Density

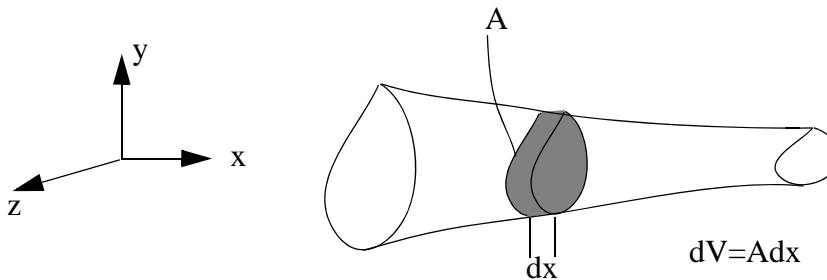
$$\text{Uniaxial tension test: } U_o = \int_0^{\varepsilon} \sigma d\varepsilon = \int_0^{\varepsilon} (E\varepsilon) d\varepsilon = \frac{E\varepsilon^2}{2} = \frac{1}{2}\sigma\varepsilon$$

$$U_o = \frac{1}{2}\tau\gamma$$

- Strain energy and strain energy density is a scalar quantity.

$$U_o = \frac{1}{2}[\sigma_{xx}\varepsilon_{xx} + \sigma_{yy}\varepsilon_{yy} + \sigma_{zz}\varepsilon_{zz} + \tau_{xy}\gamma_{xy} + \tau_{yz}\gamma_{yz} + \tau_{zx}\gamma_{zx}]$$

1-D Structural Elements



Axial strain energy

- All stress components except σ_{xx} are zero.

$$\sigma_{xx} = E\varepsilon_{xx} \quad \varepsilon_{xx} = \frac{du}{dx}(x)$$

$$U_A = \int_V \frac{1}{2} E \varepsilon_{xx}^2 dV = \int_L \left[\int_A \frac{1}{2} E \left(\frac{du}{dx} \right)^2 dA \right] dx = \int_L \left[\frac{1}{2} \left(\frac{du}{dx} \right)^2 \int_A E dA \right] dx$$

$$U_A = \int_L U_a dx \quad U_a = \frac{1}{2} EA \left(\frac{du}{dx} \right)^2$$

- U_a is the strain energy per unit length.

$$\bar{U}_A = \int_L \bar{U}_a dx \quad \bar{U}_a = \frac{1}{2} \frac{N^2}{EA}$$

Torsional strain energy

- All stress components except $\tau_{x\theta}$ in polar coordinate are zero

$$\tau_{x\theta} = G\gamma_{x\theta} \quad \gamma_{x\theta} = \rho \frac{d\phi}{dx}(x)$$

$$U_T = \int_V \frac{1}{2} G \gamma_{x\theta}^2 dV = \int_L \left[\int_A \frac{1}{2} G \left(\rho \frac{d\phi}{dx} \right)^2 dA \right] dx = \int_L \left[\frac{1}{2} \left(\frac{d\phi}{dx} \right)^2 \int_A G \rho^2 dA \right] dx$$

$$U_T = \int_L U_t dx \quad U_t = \frac{1}{2} GJ \left(\frac{d\phi}{dx} \right)^2$$

- U_t is the strain energy per unit length.

$$\bar{U}_T = \int_L \bar{U}_t dx \quad \bar{U}_t = \frac{1}{2} \frac{T^2}{GJ}$$

Strain energy in symmetric bending about z-axis

There are two non-zero stress components, σ_{xx} and τ_{xy} .

$$\sigma_{xx} = E\varepsilon_{xx} \quad \varepsilon_{xx} = -y \frac{d^2 v}{dx^2}$$

$$U_B = \int_V \frac{1}{2} E \varepsilon_{xx}^2 dV = \int_L \left[\int_A \frac{1}{2} E \left(y \frac{d^2 v}{dx^2} \right)^2 dA \right] dx = \int_L \left[\frac{1}{2} \left(\frac{d^2 v}{dx^2} \right)^2 \int_A E y^2 dA \right] dx$$

$$U_B = \int_L U_b dx \quad U_b = \frac{1}{2} EI_{zz} \left(\frac{d^2 v}{dx^2} \right)^2$$

- where U_b is the bending strain energy per unit length.

$$\bar{U}_B = \int_L \bar{U}_b dx \quad \bar{U}_b = \frac{1}{2} \frac{M_z^2}{EI_{zz}}$$

The strain energy due to shear in bending is: $U_S = \int_V \frac{1}{2} \tau_{xy} \gamma_{xy} dV = \int_V \frac{1}{2} \frac{\tau_{xy}^2}{E} dV$

As $\tau_{max} \ll \sigma_{max}$

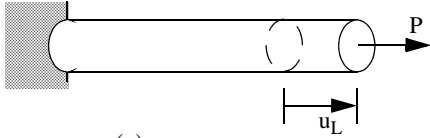
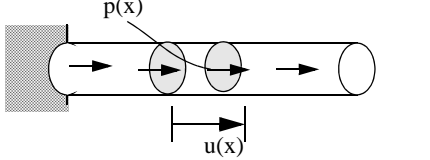
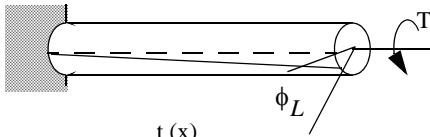
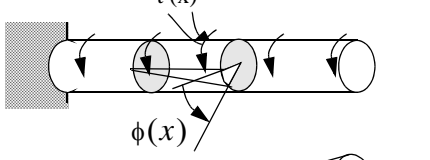
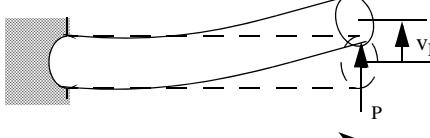
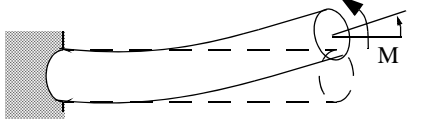
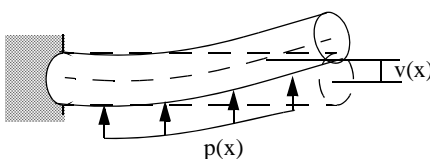
$$U_S \ll U_B$$

Work

- If a force moves through a distance, then work has been done by the force.

$$dW = F du$$

- Work done by a force is conservative if it is path independent.
- Non-linear systems and non-conservative systems are two independent description of a system.

Loading Mode	Work
	$\delta W = P \delta u_L$
	$\delta W = \int_0^L p(x) \delta u(x) dx$
	$\delta W = T \delta \phi_L$
	$\delta W = \int_0^L t(x) \delta \phi(x) dx$
	$\delta W = P \delta v_L$
	$\delta W = M \delta \theta_L$
	$\delta W = \int_0^L p(x) \delta v(x) dx$

- Any variable that can be used for describing deformation is called the generalized displacement.
- Any variable that can be used for describing the cause that produces deformation is called the generalized force.

Virtual Work

- Virtual work methods are applicable to linear and non-linear systems, to conservative as well as non-conservative systems.

The principle of virtual work:

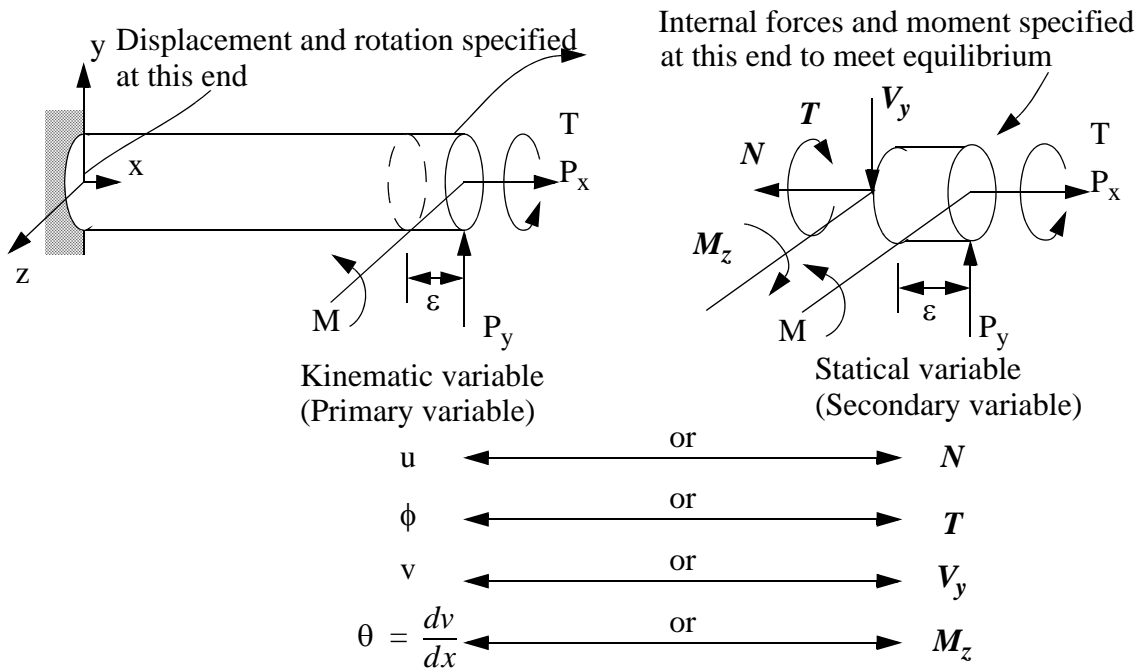
The total virtual work done on a body at equilibrium is zero.

$$\delta W = 0$$

- Symbol δ will be used to designate a virtual quantity

$$\delta W_{ext} = \delta W_{int}$$

Types of boundary conditions



Geometric boundary conditions (Kinematic boundary conditions)
(Essential boundary conditions):

Condition specified on kinematic (primary) variable at the boundary.

Statical boundary conditions
(Natural boundary conditions)

Condition specified on statical (secondary) variable at the boundary.

Kinematically admissible functions

- Functions that are continuous and satisfies all the kinematic boundary conditions are called *kinematically admissible functions*.
- actual displacement solution is always a kinematically admissible function
- Kinematically admissible functions are not required to correspond to solutions that satisfy equilibrium equations.

Statically admissible functions

- Functions that satisfy satisfies all the static boundary conditions, satisfy equilibrium equations at all points, and are continuous at all points except where a concentrated force or moment is applied are called *statically admissible functions*.
- Actual internal forces and moments are always statically admissible.
- Statically admissible functions are not required to correspond to solutions that satisfy compatibility equations.

7.3 Determine a class of kinematically admissible displacement functions for the beam shown in Figure P7.3.

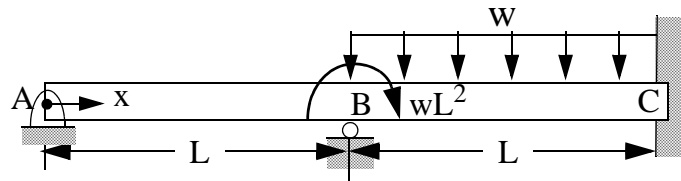
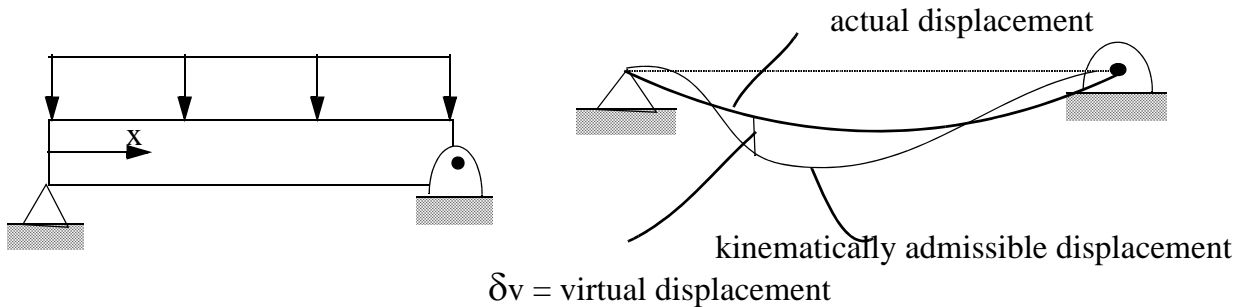


Figure P7.3

7.4 For the beam and loading shown in Figure P7.3 determine a statically admissible bending moment.

Virtual displacement method

- The virtual displacement is an infinitesimal imaginary kinematically admissible displacement field imposed on a body.



- Of all the virtual displacements the one that satisfies the virtual work principle is the actual displacement field.

Virtual Force Method

- The virtual force is an infinitesimal imaginary statically admissible force field imposed on a body.
- Of all the virtual force fields the one that satisfies the virtual work principle is the actual force field.

7.7 The roller at P shown in Figure P7.7 slides in the slot due to the force $F = 20\text{kN}$. Both bars have a cross-sectional area of $A = 100\text{ mm}^2$ and a modulus of elasticity $E = 200\text{ GPa}$. Bar AP and BP have lengths of $L_{AP} = 200\text{ mm}$ and $L_{BP} = 250\text{ mm}$ respectively. Determine the axial stress in the member AP by virtual displacement method.

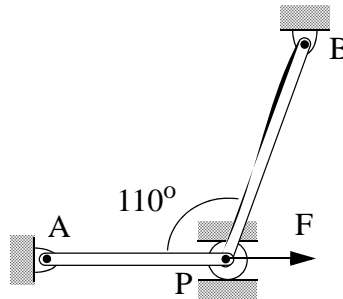


Figure P7.7

7.8 A force $F = 20\text{kN}$ is applied to pin shown in Figure P7.8. Both bars have a cross-sectional area of $A = 100\text{ mm}^2$ and a modulus of elasticity $E = 200\text{ GPa}$. Bar AP and BP have lengths of $L_{AP} = 200\text{ mm}$ and $L_{BP} = 250\text{ mm}$ respectively. Using virtual force method determine the movement of pin in the direction of force F .

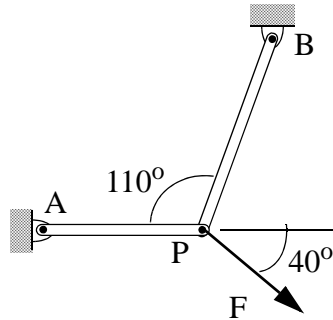


Figure P7.8

Dummy unit load method

- This is a virtual force method that is formalized.
- Can be used for axial, torsion or bending problems.

Application to beam bending

Displacement Calculations

Consider two beams.

BEAM 1: Actual beam with actual internal moment $M_1(x)$ and actual displacement $v_1(x)$.

BEAM 2: A beam with same supports as beam 1 with a *unit force* placed at point x_p at which we want to calculate the displacement. $M_2(x)$ be the statically admissible bending moment and $v_2(x)$ be the kinematically admissible displacement for beam 2.

Note: No relationship between M_2 and v_2

The internal and external virtual work for beam 2:

$$\delta W_{int} = \int_0^L M_2(x) d\theta_2 = \int_0^L M_2(x) \frac{d\theta_2}{dx} dx = \int_0^L M_2(x) \frac{d}{dx} \left(\frac{dv_2}{dx} \right) dx = \int_0^L M_2(x) \frac{d^2 v_2}{dx^2} dx$$

$$\delta W_{ext} = (1)v_2(x_p)$$

By theorem of virtual work: $v_2(x_p) = \int_0^L M_2(x) \frac{d^2 v_2}{dx^2} dx$

$v_1(x)$ is a kinematically admissible displacement field, hence can be used for $v_2(x)$.

$$v_1(x_p) = \int_0^L M_2(x) \frac{d^2 v_1}{dx^2} dx = \int_0^L \frac{M_2(x) M_1(x)}{EI} dx$$

Slope Calculations

BEAM 1: Actual beam with actual internal moment $M_1(x)$ and actual displacement $v_1(x)$.

BEAM 2: A beam with same supports as beam 1 with a *unit moment* placed at point x_p at which we want to calculate the slope. $M_2(x)$ be the statically admissible bending moment and $v_2(x)$ be the kinematically admissible displacement for beam 2.

Note: No relationship between M_2 and v_2

The internal and external virtual work for beam 2:

$$\delta W_{int} = \int_0^L M_2(x) d\theta_2 = \int_0^L M_2(x) \frac{d\theta_2}{dx} dx = \int_0^L M_2(x) \frac{d}{dx} \left(\frac{dv_2}{dx} \right) dx = \int_0^L M_2(x) \frac{d^2 v_2}{dx^2} dx$$

$$\delta W_{ext} = (1) \frac{dv_2}{dx}(x_p)$$

By theorem of virtual work: $\frac{dv_2}{dx}(x_p) = \int_0^L M_2(x) \frac{d^2 v_2}{dx^2} dx$

$v_1(x)$ is a kinematically admissible displacement field, hence can be used for $v_2(x)$.

$$\frac{dv_1}{dx}(x_p) = \int_0^L M_2(x) \frac{d^2 v_1}{dx^2} dx = \int_0^L \frac{M_2(x) M_1(x)}{EI} dx$$

7.21 Using dummy unit load method, find the reaction force at A and deflection at B in terms of P , E , I , and L for the beam shown in Figure P7.21.

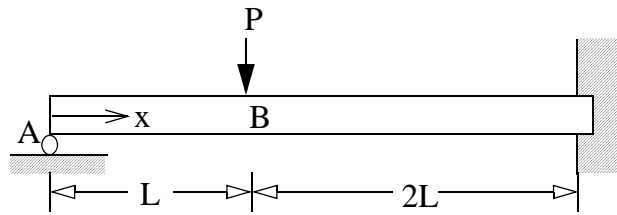


Figure P7.21

Castigliano's theorem

- Simple and more elegant way of finding reaction forces and/or moments for statically indeterminate structures.

Instead of a unit force we consider a force F applied at x_p in the dummy unit load method. For linear systems the corresponding statically admissible moment (\tilde{M}_2) would be F multiplied by M_2 .

$$\tilde{M}_2 = FM_2 \quad \text{or} \quad M_2 = \frac{\partial \tilde{M}_2}{\partial F} \quad v_1(x_p) = \int_0^L \frac{M_2(x)M_1(x)}{EI} dx = \int_0^L \frac{1}{EI} \left(\frac{\partial \tilde{M}_2}{\partial F} M_1(x) \right) dx$$

The actual moment is a statically admissible moment, and hence we can substitute $\tilde{M}_2 = M_1$ we obtain the following:

$$v_1(x_p) = \int_0^L \frac{1}{EI} \left(\frac{\partial M_1}{\partial F} M_1(x) \right) dx = \int_0^L \frac{1}{2EI} \left(\frac{\partial M_1^2}{\partial F} \right) dx = \frac{\partial}{\partial F} \left[\int_0^L \frac{M_1^2}{2EI} dx \right] = \frac{\partial \bar{U}_B}{\partial F}$$

- The derivative of the complimentary strain energy with respect to a force at x_p gives the deflection in the direction of the force at x_p .

$$\boxed{\frac{dv_1}{dx}(x_p) = \frac{\partial \bar{U}_B}{\partial M}}$$

- The derivative of the complimentary strain energy with respect to a moment at x_p gives the slope in the direction of the moment at x_p .
- Performing the derivative with respect to force and moment before performing integration will generally result in less algebra.
- The integrals obtained after taking the derivative with respect to force and moment result in integrals that are identical to the dummy unit load method for finding reactions.

7.27 Using Castigliano's theorem, find the reaction force at A and deflection at B in terms of P , E, I , and L for the beam shown in Figure P7.21.

